A model for the emergence of cooperation, interdependence and structure in evolving networks

Sanjay Jain[*†‡1] and **Sandeep Krishna**[§2]

[*] Centre for Theoretical Studies and [§] Physics Department, Indian Institute of Science,
Bangalore 560 012, India
[†] Santa Fe Institute, 1399 Hyde Park Road, Santa Fe, NM 87501, USA
[‡] Jawaharlal Nehru Centre for Advanced Scientific Research, Bangalore 560 064, India
[1] *jain@santafe.edu, jain@cts.iisc.ernet.in*; [2] *sandeep@physics.iisc.ernet.in*

Evolution produces complex and structured networks of interacting components in chemical, biological, and social systems. We describe a simple mathematical model for the evolution of an idealized chemical system to study how a network of cooperative molecular species arises and evolves to become more complex and structured. The network is modeled by a directed weighted graph whose positive and negative links represent 'catalytic' and 'inhibitory' interactions among the molecular species, and which evolves as the least populated species (typically those that go extinct) are replaced by new ones. A small autocatalytic set (ACS), appearing by chance, provides the seed for the spontaneous growth of connectivity and cooperation in the graph. A highly structured chemical organization arises inevitably as the ACS enlarges and percolates through the network in a short, analytically determined time scale. This self-organization does not require the presence of self-replicating species. The network also exhibits catastrophes over long time scales triggered by the chance elimination of 'keystone' species, followed by recoveries.

Structured networks of interacting components are a hallmark of several complex systems, for example, the chemical network of molecular species in cells [1], the web of interdependent biological species in ecosystems [2,3], and social and economic networks of interacting agents in societies [4–7]. The structure of these networks is a product of evolution, shaped partly by the environment and physical constraints and partly by the population (or other) dynamics in the system. For example, imagine a pond on the prebiotic earth containing a set of interacting molecular species with some concentrations. The interactions among the species in the pond affect how the populations evolve with time. If some population goes to zero, or if new molecular species enter the pond from the environment (through storms, floods or tides), the effective chemical network existing in the pond changes. We discuss a mathematical model that attempts to incorporate this interplay between a network, populations, and the environment in a simple and idealized fashion. The model (including an earlier version [8,9]) was inspired by the ideas and results in refs. [10–18]. Related but different models are studied in refs. [19–21].

The Model

The system consists of s species labeled by the index $i = 1, 2, \ldots, s$. The network of interactions between species is specified by the $s \times s$ real matrix $C \equiv \{c_{ij}\}$. The network can be visualised as a directed graph whose nodes represent the species. A non-zero c_{ij} is represented by a directed weighted link from node j to node i. If $c_{ij} > 0$ then the corresponding link is a cooperative link: species j catalyzes the production of species i. If $c_{ij} < 0$ it is a destructive link: the presence of j causes a depletion of i.

Population dynamics. The model contains another dynamical variable $\mathbf{x} \equiv (x_1, \ldots, x_s)$, where x_i stands for the relative population of the ith species ($0 \leq x_i \leq 1$, $\sum_{i=1}^{s} x_i = 1$). The time evolution of $\mathbf{x}$ depends upon the interaction coefficients C, as is usual in population models. The specific evolution rule we consider is

$$
\begin{aligned}
\dot{x}_i &= f_i &&\text{if } x_i > 0 \quad \text{or} \quad f_i \geq 0, \\
&= 0 &&\text{if } x_i = 0 \quad \text{and} \quad f_i < 0,
\end{aligned}
\tag{1}
$$
$$
\text{where } f_i = \sum_{j=1}^{s} c_{ij} x_j - x_i \sum_{k,j=1}^{s} c_{kj} x_j.
$$

This is a particularly simple idealization of catalyzed chemical reaction dynamics in a well stirred reactor, as explained later.

Graph dynamics. The dynamics of C in turn depends upon $\mathbf{x}$, as follows: Start with a random graph of s nodes: c_{ij} is non-zero with probability p and zero with probability $1 - p$. If nonzero, c_{ij} is chosen randomly in the interval $[-1, 1]$ for $i \neq j$ and $[-1, 0]$ for $i = j$. Thus a link between two distinct species, if it exists, is just as likely to be cooperative as destructive, and a link from a species to itself can only be inhibitive, i.e., autocatalytic or self-replicating individual species are not allowed. The variable $\mathbf{x}$ is initialized by choosing each x_i randomly between 0 and 1, and then rescaling all x_i uniformly such that $\sum_{i=1}^{s} x_i = 1$. The evolution of the network proceeds in three steps :

1) Keeping the network fixed, the populations are evolved according to (1) for a time T which is large enough for $\mathbf{x}$ to get reasonably close to its attractor. We denote $X_i \equiv x_i(T)$.

2) The set of nodes i with the least value of X_i is determined. We call this the set of 'least fit' nodes, identifying the relative population of a species in the attractor (or, more specifically, at T) with its 'fitness' in the environment defined by the graph. One of the least fit nodes is chosen randomly (say i_0) and removed from the system along with all its links leaving a graph of $s - 1$ species.

3) A new node is added to the graph so that it again has s nodes. The links of the added node (c_{ii_0} and $c_{i_0 i}$, for $i = 1, \ldots, s$) are assigned randomly according to the same rule as for the nodes in the initial graph. The new species is given a small relative population $x_{i_0} = x_0$ and the other populations are rescaled to keep $\sum_{i=1}^{s} x_i = 1$. This process, from step 1 onwards, is iterated many times.

Motivation for model structure. The choice of Eq. (1) is motivated by the rate equations in a well stirred chemical reactor (representing, say, a prebiotic pond) as follows: If species j catalyses the ligation of reactants A and B to form the species i, $A + B \xrightarrow{j} i$, then the rate of growth of the population y_i of species i will be given by $\dot{y}_i = k(1 + \nu y_j)n_A n_B - \phi y_i$, where n_A, n_B are reactant concentrations, k is the rate constant for the spontaneous reaction, ν is the catalytic efficiency, and ϕ represents a common death rate or dilution flux in the reactor [22]. Assuming the catalysed reaction is much faster than the spontaneous reaction, and the concentrations of the reactants are large and fixed, the rate equation becomes $\dot{y}_i = c y_j - \phi y_i$, where c is a constant. If species i has multiple catalysts, we get $\dot{y}_i = \sum_{j=1}^{s} c_{ij} y_j - \phi y_i$.

The first of equations (1) follows from this upon using the definition $x_i = y_i / \sum_{j=1}^{s} y_j$. When negative links are permitted, the second of equations (1) is needed in general to prevent relative populations from going negative. (With negative links, a more realistic chemical interpretation would be obtained if $\dot{x}_i$ were proportional to x_i, but for simplicity we retain the form of (1) in this paper.) (1) may be viewed as defining an artificial chemistry in the spirit of refs. [13–17].

The rules for the evolution of the network C are intended to capture two key features of natural evolution, namely, selection and novelty. The species that has the least population in the attractor configuration is the one most likely to be eliminated in a large fluctuation in a possible hostile environment. Often, the least value of X_i is zero. Thus the model implements selection by eliminating from the network a species that has become extinct or has the least chance of survival [18]. Novelty is introduced in the network in the form of a new species. This species has on average the same connectivity as the initial set of species, but its actual connections with the existing set are drawn randomly. E.g., if a storm brings into a prebiotic pond a new molecular species from the environment, the new species might be statistically similar to the one being eliminated, but its actual catalytic and inhibitory interactions with the surviving species can be quite different. Another common feature of natural evolution is that populations typically evolve on a fast time scale compared to the network. This is captured in the model by having the x_i relax to their attractor before the network is updated. The idealization of a fixed total number of species s is one that we hope to relax in future work.

The model described above differs from the one studied in [8, 9] in that it allows negative links and varying link strengths, and that the population dynamics, given by (1), is no longer linear. The earlier model had only fixed point attractors; here limit cycles are also observed. Since C now has negative entries, the formalism of non-negative matrices no longer applies.

Results

Emergence of cooperation and interdependence. Figure 1 shows a sample run. The same qualitative behaviour is seen in each of several hundred runs performed for p values ranging from 0.00002 to 0.01 and for $s = 100, 150, 200$. The fact that the ratio of number of cooperative to destructive links

at first remains constant at unity (statistically) and then increases by more than an order of magnitude is evidence of the emergence of cooperation. The increase in $\bar{d}$ by an order of magnitude is a quantitative measure of the increase of interdependence of species in the network. The increase in the total density of links $(l_+ + l_-)/s$ is another aspect of the increase of complexity of the system. Note that in the model selection rewards only 'performance' as measured in terms of relative population; the rules do not select for higher cooperativity per se. Since a new species is equally likely to have positive or negative links with other species, the introduction of novelty is also not biased in favour of cooperativity. The fact that this behaviour is not a consequence of any intrinsic bias in the model that favours the increase of cooperation and interdependence is evidenced by the flat initial region of all the curves.

Autocatalytic sets. The explanation for the above behaviour lies in the formation and growth of certain structures, autocatalytic sets (ACSs), in the graph. An ACS is defined as a set of nodes such that each node has at least one incoming positive link from a node in the set. Thus an ACS has the property of catalytic closure, i.e., it contains a catalyst for each of its members [23–25]. The simplest example of an ACS is a cycle of positive links. Every ACS is not such a cycle but it can be shown that an ACS must contain a cycle of positive links [9]. In Fig. 1, there is no ACS in the graph until $n = 1903$. A small ACS (which happens to be a cycle of positive links between two nodes) appears at $n \equiv n_1 = 1904$, exactly where the behaviour of the s_1 curve changes. As time proceeds this ACS becomes more complex and enlarges until at $n \equiv n_2 = 3643$ the entire graph becomes an ACS. l_+ and $\bar{d}$ exhibit an increase and l_- a decrease as the ACS comes to occupy a significant part of the graph. After the ACS first appears (at $n = n_1$), the set of populated nodes in the attractor configuration (s_1 in number), is always an ACS (except for certain catastrophic events to be discussed later), which we call the 'dominant ACS'. The spontaneous appearance of a small ACS at some $n = n_1$, its persistence (except for catastrophes), and its growth until it spans the graph at $n = n_2$, is seen in each of the several hundred runs mentioned earlier. The growth of the ACS across the graph between n_1 and n_2 occurs exponentially (with stochastic fluctuations),

$$s_1(n) \approx s_1(n_1)e^{(n-n_1)/\tau_g}, \qquad \tau_g = 2/p. \quad (2)$$

This agrees with simulations as shown in Figure 2. The average time scale $\tau_a \equiv \langle n_1 \rangle$ for the first appearance of the ACS is given, for sufficiently small p, by $\tau_a \approx 4/(p^2 s)$ (= 1600 for $p = 0.005$ and $s = 100$).

Upto $n = n_1$, the graph has no ACS. It has chains and trees of positive and negative links and possibly loops containing negative links. These latter structures are not robust. For example, consider a chain of two positive links $1 \rightarrow 2 \rightarrow 3$. Since catalytic links are pointing to node 3, it will do well populationally compared to nodes 1 and 2. However, since 1 has no incoming catalytic links, its relative population will decline to zero under (1), and it can be picked for replacement in the next graph update. This can disrupt the chain and hence erode the 'well being' of node 3 until eventually after some graph updates the latter can also join the ranks of the least fit. Species 3 gets eliminated eventually because it does not feedback into and 'protect' species 1 and 2, on whom its 'well being' depends. In a graph without an ACS no structure is protected from disruption. Since every node is liable to be replaced sooner or later, the graph remains as random as the initial graph (we have checked that the probability distribution of the number of incoming and outgoing links at a node remains the appropriate binomial for $n < n_1$). This explains why s_1, $l_\pm$ and $\bar{d}$ hover around their initial values.

The picture changes the moment a small ACS appears in the graph. The key point is that by virtue of catalytic closure, members of the ACS do well *collectively* in the population dynamics governed by (1). An ACS is a collective self-replicator and beats chains, trees and other non ACS structures in the population game, reducing their X_i to zero when it appears. Thus, since graph update proceeds by replacing one of the nodes with $X_i = 0$ (if present) with a new one, such a replacement being outside the dominant ACS can cause no damage to the links that constitute the ACS. That is why the ACS structure, once it appears, is much more robust than the non-ACS structures discussed earlier. If the new node happens to get an incoming positive link from the dominant ACS, it becomes part of it. Thus the dominant ACS tends to expand in the graph as new nodes get attached to it [15] [8,9] and s_1 increases. In Δn graph updates the average increase in s_1, which is the number of added nodes which will get a positive link from one of the s_1 nodes of the dominant ACS, is $\Delta s_1 \approx (p/2)s_1 \Delta n$, for small p. This proves (2). (Note that the exponential growth described by (2) is not to be confused with the exponential growth of

3

populations y_i of species that are part of the ACS. (2) reflects the growth of the ACS across the graph, or the increase in the number of species that constitute the ACS.)

Since the dominant ACS grows by adding positive links from the existing dominant ACS, the number of positive links increases as the ACS grows. On the other hand nodes receiving negative links usually end up being least fit, hence negative links get removed when these nodes are eliminated. Which novelty is captured thus depends upon the existing 'context'; the network evolves by preferentially capturing links and nodes that 'latch on' cooperatively to the existing ACS and disregarding those that do not. The 'context' itself arises when the ACS structure first appears; this event transforms the nature of network evolution from random to 'purposeful' (in this case directed towards increasing cooperation). Before the ACS appears nothing interesting happens even though selection is operative (the least populated species are being eliminated). It is only after the ACS topological structure appears that selection for cooperation and complexity begins.

Inevitability of autocatalytic sets. Note that the appearance of an ACS, though a chance event, is inevitable. For $sp \ll 1$, the probability that a graph not containing a two cycle will acquire one at the next time step is $p^2 s/4 \equiv q$. Since the probability of occurrence of 3-cycles, etc., is much smaller, the probability distribution of arrival times n_1 is approximated by $P(n_1) = q(1-q)^{n_1-1}$, whose mean τ_a is $1/q$. Since this probability declines exponentially after a time scale $1/q$, the appearance of an ACS is inevitable, even for arbitrarily small (but finite) p.

Occasionally in a graph update s_1 can decrease for various reasons. If the new node forms an ACS of its own with nodes outside the dominant ACS, and the new ACS has a higher population growth rate (as determined by (1)) than the old ACS it drives the species of the latter to extinction and becomes the new dominant ACS. Alternatively the new node could be a 'destructive parasite': it receives one or more positive links from and gives one or more negative links to the dominant ACS. Then part or whole of the ACS may join the set of least fit nodes. Structures that diminish the size of the dominant ACS or destroy it appear rarely. For example in Figure 1, destructive parasites appeared 6 times at $n = 3388, 3478, 3576, 3579, 3592$ and 3613. In each case s_1 decreased by 1.

Emergence of structure. At $n = n_2$ the whole graph becomes an ACS; the entire system can collectively self-replicate despite the explicit absence of individual self-replicators. Such a fully autocatalytic set is a very non-random structure. Consider a graph of s nodes and let the probability of a positive link existing between any pair of nodes be p^*. Such a graph has on average $m^* = p^*(s-1)$ incoming or outgoing positive links per node. For the entire graph to be an ACS, each node must have at least one incoming positive link, i.e., each row of the matrix C must contain at least one positive element. Hence the probability, P, for the entire random graph to be an ACS is

$$
\begin{aligned}
P &= \text{probability that every row has at least} \\
&\qquad \text{one positive entry} \\
&= [\text{probability that a row has at least} \\
&\qquad \text{one positive entry}]^s \\
&= [1 - (\text{probability that every entry} \\
&\qquad \text{of a row is} \le 0)]^s \\
&= [1 - (1-p^*)^{s-1}]^s \\
&= [1 - (1 - m^*/(s-1))^{s-1}]^s.
\end{aligned}
$$

For large s and $m^* \sim O(1)$,

$$P \approx (1 - e^{-m^*})^s = e^{-\alpha s}, \qquad (3)$$

where α is positive and $O(1)$. At $n = n_2$, we find in all our runs that $l_+(n_2) \equiv l^*$ is greater than s but of order s, i.e., $m^* \sim O(1)$. Thus dynamical evolution in the model via the ACS mechanism converts a random organization into a highly structured one that is exponentially unlikely to appear by chance. In the displayed run at $n = n_2$ the graph had 117 positive links. The probability that a random graph with $s = 100$ nodes and $m^* = 1.17$ would be an ACS is given by (3) to be $\approx 10^{-16}$.

Such a structure would take an exponentially long time to arise by pure chance. The reason it arises inevitably in a short time scale in the present model is the following: a *small* ACS can appear by chance quite readily, and once appeared, it grows exponentially fast across the graph by the mechanism outlined earlier. The dynamical appearance of such a structure may be regarded as the emergence of 'organizational order'. The appearance of 'exponentially unlikely' structures in the prebiotic context has been a puzzle. The fact that in the present model such structures inevitably form in a short time may be relevant for the origin of life problem.

The self-organization time scale in a prebiotic scenario. We now speculate on a possible application to prebiotic chemical evolution. Imagine the

molecular species to be small peptide chains with weak catalytic activity in a prebiotic pond alluded to earlier. The pond periodically receives an influx of new molecular species being randomly generated elsewhere in the environment, through tides, storms or floods. Between these influxes of novelty the pond behaves as a well stirred reactor where the populations of existent molecular species evolve according to (a realistic version of) eq. (1) and reach their attractor configuration. Under the assumption that the present model captures what happens in such a pond, the growth timescale (2) for a highly structured almost fully autocatalytic chemical organization in the pond is $\tau_g = 2/p$ in units of the graph update time step. In this scenario, the latter time unit corresponds to the periodicity of the influx of new molecular species, hence it ranges from one day (for tides) to one year (for floods). Further, in the present model $p/2$ is the probability that a random small peptide will catalyse the production of another [25], and this has been estimated in [12]: $p/2 \sim 10^{-5} - 10^{-10}$. With $p/2 \sim 10^{-8}$, for example, the time scale for a highly structured chemical organization to grow in the pond would be estimated to be of the order of 10^6 to 10^8 years. It is believed that life originated on Earth in a few hundred million years after the oceans condensed. These considerations suggest that it might be worthwhile to empirically pin down the 'catalytic probability' p (introduced in [25]) for peptides, catalytic RNA, lipids, etc., on the one hand, and explore chemically more realistic models on the other.

Catastrophes and recoveries in the network dynamics. After $n = n_2$ the character of the network evolution changes again. For the first time the least fit node will be one of the ACS members. Most of the time elimination of the node does not affect the ACS significantly and s_1 fluctuates between s and $s - 1$. Sometimes the least fit node could be a 'keystone' species which plays an important organizational role in the network despite its low population. When such a node is eliminated many other nodes can get disconnected from the ACS resulting in large dips in s_1 and $\bar{d}$ and subsequently large fluctuations in l_+ and l_-. These large 'extinction events' can be seen in Figure 3. Occasionally the ACS can even be destroyed completely. The system recovers on the time scale τ_g after large extinctions if the ACS is not completely destroyed; if it is and the next few updates obliterate the memory of previous structures in the graph, then again a time on average τ_a elapses before

an an ACS arises and the self-organization process begins anew. It may be of interest (especially in ecology, economics, and finance), that network dynamics based on a fitness selection and the 'incremental' introduction of novelty, as discussed here, can by itself cause catastrophic events without the presence of large external perturbations.

Discussion

We have described an evolutionary model in which the dynamics of species' populations (fast variables) and the graph of interactions among them (slow variables) are mutually coupled. The network dynamics displays self-organization seeded by the chance but inevitable appearance of a small cooperative structure, namely, an ACS. In a dynamics that penalizes species for low population performance, the collective cooperativity of the ACS members makes the set relatively robust against disruption. New species that happen to latch on cooperatively to this structure preferentially survive, further enlarging the ACS in the process. Eventually the graph acquires a highly non-random structure. We have discussed the time evolution of quantitative measures of cooperation, interdependence and structure of the network, which capture various aspects of the complexity of the system.

It is noteworthy that collectively replicating ACSs arise even though individual species are not self-replicating. Thus the present mechanism is different from the hypercycle [26], where a template is needed to produce copies of existing species. Unlike the hypercycle, the ACS is not disrupted by parasites and short-circuits and grows in complexity, as evidenced in all our runs. It can be disrupted, however, when it loses a 'keystone' species.

It is also worth mentioning one departure from [12], in that we find that a fully autocatalytic system (or percolating ACS) is not needed apriori for self-organization. In the present model a small ACS, once formed, typically expands (see also [15]) and eventually percolates the whole network dynamically. This dynamical process might be relevant for economic takeoff and technological growth in societies.

References

[1] Lodish, H., Baltimore, D., Berk, A., Zipursky, S. L., Matsudaira, P. & Darnell, J. E. (1995) *Molecular Cell Biology* (Scientific American Books Inc., New York).

[2] Cohen, J. E., Briand, F. & Newman, C. M. (1990) *Community Food Webs: Data and Theory* (Springer Verlag, New York).

[3] Pimm, S. L. (1991) *The Balance of Nature* (Univ. Chicago Press).

[4] Wellman, B. & Berkovitz, S. D., eds. (1988) *Social Structures: A Network Approach* (Cambridge Univ. Press).

[5] Axelrod, R. (1997) *The Complexity of Cooperation* (Princeton Univ. Press).

[6] Watts, D.J. & Strogatz, S. H. (1998) *Nature* **393**, 440-442.

[7] Barabasi, A.-L. & Albert, R. (1999) *Science* **286**, 509-512.

[8] Jain, S. & Krishna, S. (1998) *Phys. Rev. Lett.* **81**, 5684-5687 (*adap-org/9809003*).

[9] Jain, S. & Krishna, S. (1999) *Computer Physics Comm.* **121-122**, 116-121 (*adap-org/9810005*).

[10] Dyson, F. (1985) *Origins of Life* (Cambridge Univ. Press).

[11] Kauffman, S. A. (1986) *J. Theor. Biol.* **119**, 1-24.

[12] Kauffman, S. A. (1993) *The Origins of Order* (Oxford Univ. Press).

[13] Farmer, J. D., Kauffman, S. A., & Packard, N. H. (1986) *Physica* **D22**, 50-67.

[14] Bagley, R. J. & Farmer, J. D. (1991) in *Artificial Life II*, eds. Langton, C. G., Taylor, C., Farmer, J. D. & Rasmussen, S. (Addison Wesley, Redwood City), pp. 93-140.

[15] Bagley, R. J., Farmer, J. D. & Fontana, W. (1991) in *Artificial Life II*, eds. Langton, C. G., Taylor, C., Farmer, J. D. & Rasmussen, S. (Addison Wesley, Redwood City), pp. 141-158.

[16] Fontana, W. (1991) in *Artificial Life II*, eds. Langton, C. G., Taylor, C., Farmer, J. D. & Rasmussen, S. (Addison Wesley, Redwood City), pp. 159-209.

[17] Fontana, W. & Buss, L. (1994) *Bull. Math. Biol.* **56**, 1-64.

[18] Bak, P. & Sneppen, K. (1993) *Phys. Rev. Lett.* **71**, 4083-4086.

[19] Happel, R. & Stadler, P. F. (1998) *J. Theor. Biol.* **195**, 329-338.

[20] Yasutomi, A. & Tokita, K. (1998) Preprint available from http://paradox.harvard.edu/∼ tokita/list.html .

[21] Segré, D., Ben-Eli, D. & Lancet, D. (2000) *Proc. Natl. Acad. Sc. USA* **97**, 4112-4117.

[22] Ashmore, P. G. (1963) *Catalysis and Inhibition of Chemical Reactions* (Butterworth & Co., London).

[23] Eigen, M. (1971) *Naturwissenschaften* **58**, 465-523.

[24] Rossler, O. E. (1971) *Z. Naturforschung* **26b**, 741-746.

[25] Kauffman, S.A. (1971) *J. Cybernetics* **1**, 71-96.

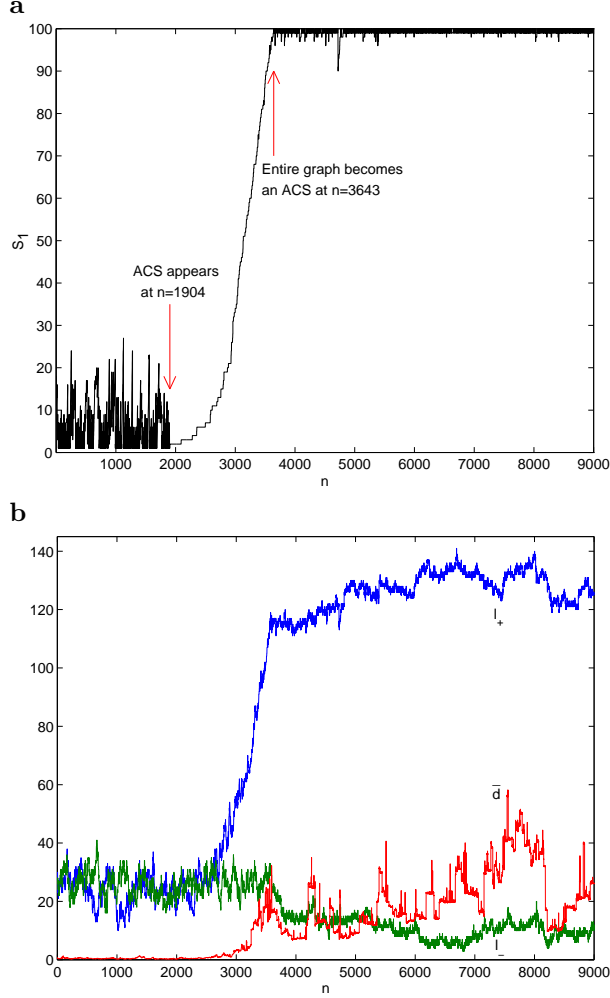[26] Eigen, M. & Schuster, P. (1979) *The Hypercycle* (Springer, Berlin).

a

b

Figure 1. A run with parameter values $s = 100$, $p = 0.005$, and $x_0 = 10^{-4}$.
(**a**) Number of populated species, s_1, in the attractor of (1) (i.e., number of nodes with $X_i > 0$) after the nth addition of a new species (i.e., after n graph update time steps). (**b**) The number, l_+, of positive links ($c_{ij} > 0$) in the graph (blue); the number, l_-, of negative links (green); and 'interdependency', $\bar{d}$, of the species in the network (red). Interdependency is defined as $\bar{d} \equiv (1/s) \sum_{i=1}^{s} d_i$, where d_i is the 'dependency' of the ith node. $d_i \equiv \sum_{kj} |c_{kj}| h_k^i$, where h_k^i is 1 if there exists a directed path from k to i and 0 otherwise. d_i is the sum of (the absolute value of) the strengths of all links that eventually feed into i along some directed path. The curves have three distinct regions. Initially s_1 is small; most of the species have zero relative populations. l_+ and l_- also do not vary much from their initial (random graph) value ($\approx p s^2/2 = 25$) and remain approximately equal. $\bar{d}$ hovers about its initial low value. In the second region s_1, l_+ and $\bar{d}$ show a sharp increase and l_- decreases. In the third region s_1, $l_\pm$ and $\bar{d}$ level off (but with fluctuations) and almost all species have non-zero populations in contrast to the initial period.
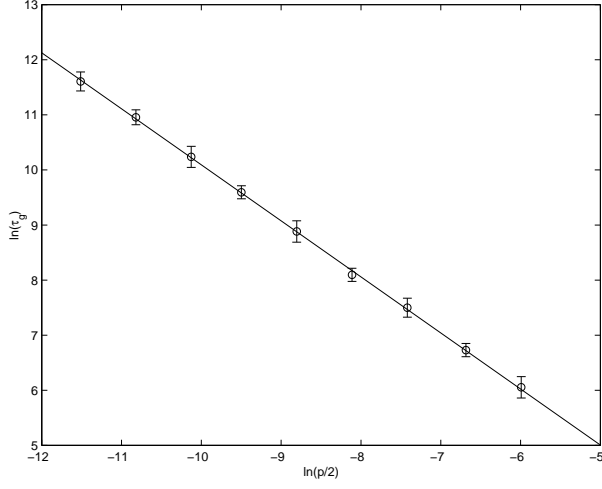
Figure 2. Power law dependence of τ_g on p. Each data point shows the average of τ_g over 5 different runs with $s = 100$ and the given p value. The error bars correspond to one standard deviation. The best fit line has slope -1.02 ± 0.03 and intercept -0.08 ± 0.26 which is consistent with the expected slope -1 and intercept 0.
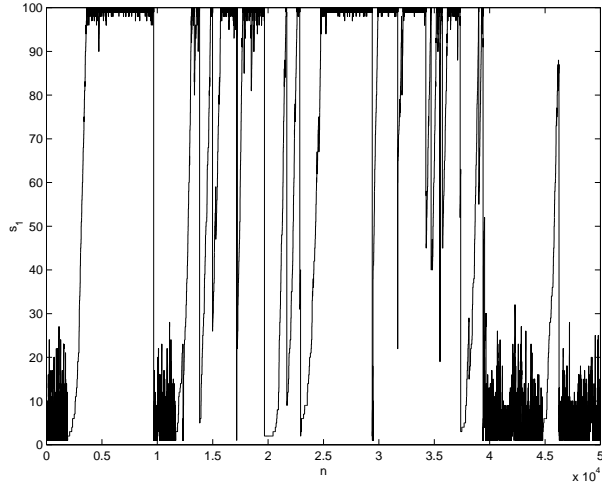

Figure 3. The run of Figure 1a displayed over a much longer time scale.